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one of the major subjects in studies of high  $T_c$ 

superconductivity in oxide materials. The impor-

tance of the BKBO system lies in the fact that some

of its superconducting properties are consistent

with the conventional s-wave isotropic supercon-

ductivity, but others resemble high- $T_c$  cuprates. In

particular, substantial isotope effect [1,2], strong

phonon spectrum from the neutron scattering

measurements [3], onset of a superconducting gap

with  $2\Delta_0/T_c = 3.5 \pm 0.3$  from the tunnelling and

optical experiments [4-7] indicate a significant role

## Temperature dependence of magnetic penetration depth in $B_{1-x}K_xBiO_3$ superconductor

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## **Abstract**

Temperature behavior of the magnetic penetration depth,  $\lambda(T)$ , was measured in the Meissner state of  $B_{1-x}K_xBiO_3$  (x=0.37) single crystal using tunnel diode oscillator technique. Measurements were carried out down to 0.4 K. At low temperatures,  $0.013 \leqslant T/T_c \leqslant 0.4$ ,  $\lambda(T)$  is exponentially flat, which provides a strong evidence for conventional s-wave BCS behavior.

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Keywords: Magnetic penetration depth; BKBO; Pairing symmetry

Magnetic penetration depth is an effective tool to study properties of superconductors. At low temperatures, its temperature dependence is directly related to the density of low energy quasiparticles, which in turn can be related to the anisotropy of the superconducting energy gap on the Fermi surface. For investigation of the low lying excitations and thus the anisotropy of the energy gap, the analysis is considerably less ambiguous if data are taken on high quality single crystal and temperatures well below  $T_{\rm c}/3$ . In this paper we report penetration depth measurements on  $B_{0.63}K_{0.37}BiO_3$  ( $T_{\rm c}\approx 31~{\rm K}$ ) down to 0.4 K. The investigation of the superconductivity in  $B_{1-x}K_xBiO_3$  (BKBO) system has been

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of electron–phonon interactions in superconductions to BKBO. However, the low density of states at the Fermi level with  $T_c$  as high as 30 K [2] and insulator–superconductor transition by doping [8] are common with high- $T_c$  cuprates. In contrast to high  $T_c$  cuprates which have two dimensional CuO<sub>2</sub> planes, BKBO has simple three dimensional cubic perovskite structure.

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Single crystal of  $B_{1-x}K_xBiO_3$  (x=0.37) was prepared by the electro-chemical method reported elsewhere [9]. DC magnetization was measured by using *Quantum Design* MPMS SQUID magnetometer. Zero-field cooled and field cooled temperature scans taken in external field of 10 Oe are shown in Fig. 1. Superconducting transition temperature of the sample is  $T_c \approx 31$  K and the curves show regular superconducting screening with the Meissner expulsion of about 20%, which provides an indication of relatively low pinning. Magnetization loops confirm low pinning behavior.

Penetration depth measurements were carried out using a 13 MHz tunnel diode LC resonator [10,11] mounted in a He<sup>3</sup> refrigerator. The sample was placed on the movable sapphire stage with temperature control from 0.35 to 40 K. During the measurements the sample was exposed to the small ac field  $H_{ac} \simeq 20$  mOe much less than first critical field,  $H_{c1} \approx 750$  Oe at 5 K [12]. The relative resonance frequency shift,  $\Delta f = f(T) - f(T_{\rm min})$ , is related to the change of the magnetic penetration depth via  $\Delta f = -G\Delta\lambda$ , where G is a geometrical factor that depends upon the sample shape and volume as well as the coil geometry [11]. Low temperature behavior of magnetic penetration depth is shown in Fig. 2. The frequency shift in

Fig. 2 was normalized by the value of total frequency shift,  $\Delta f_0 = f(T_{\rm c}) - f(T_{\rm min}) \approx 16307$  Hz. The solid line in the figure represents the fit to a low temperature BCS expression for an s-wave material,

$$\Delta \lambda \approx \lambda(0) \sqrt{\frac{\pi \Delta_0}{2T}} \exp\left(-\frac{\Delta_0}{T}\right).$$
 (1)

Here  $\Delta_0$  is the value of the energy gap at zero temperature [13]. The fitting range was chosen up to  $0.33T_c$  to ensure the validity of the low temperature expansion.  $2\Delta_0/T_c=3.53$  corresponds to standard weak coupling BCS value. The dash-dotted line shows the low temperature behavior of magnetic penetration depth predicted for a clean d-wave superconductor [14].

$$\Delta \lambda \approx \lambda(0) \left(1 - \frac{\ln 2}{\Delta_0} T\right).$$
 (2)

The value of  $2\Delta_0/T_c$  for d-wave case was again chosen to be 3.53 in accordance with results of tunnelling and optical experiments.

Clearly the isotropic s-wave BCS curve provides best description of the low temperature penetration depth variation indicating the isotropic nature of superconducting gap for BKBO system. Some apparent noise in the data is because in the

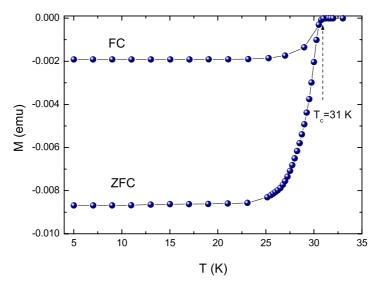


Fig. 1. Temperature dependence of magnetic moment in B–K–Bi–O single crystal in zero-field cooled and field cooled experiment in an external applied field of 10 Oe.

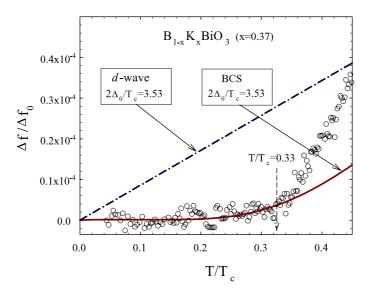


Fig. 2. Low temperature penetration depth variation in  $B_{1-x}K_xBiO_3$  (x = 0.37) single crystal. The solid line shows the low temperature exponential fit to the weak-coupling BCS expression. Dash-dotted line represents low temperature d-wave behavior. See text for details

temperature interval of interest, the penetration depth is exponentially flat with no systematic temperature dependence.

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## References

- [1] C.K. Loong, D.G. Hinks, P. Vashishta, W. Jin, R.K. Kalia, M.H. Degani, D.L. Price, J.D. Jorgensen, B. Dabrowski, A.W. Mitchell, D.R. Richards, Y. Zheng, Phys. Rev. Lett. 66 (1991) 3217.
- [2] B. Batlogg, R.J. Cava, L.W. Rupp Jr., A.M. Mujsce, J.J. Krajewski, J.P. Remeika, W.F. Peck Jr., A.S. Cooper, G.P. Espinosa, Phys. Rev. Lett. 61 (1988) 1670.
- [3] C.K. Loong, P. Vashishta, R.K. Kalia, M.H. Degani, D.L. Price, J.D. Jorgensen, D.G. Hinks, B. Dabrowski, A.W.

- Mitchell, D.R. Richards, Y. Zheng, Phys. Rev. Lett. 62 (1989) 2628.
- [4] Q. Huang, J.F. Zasadzinski, N. Tralshawala, K.E. Gray, D.G. Hinks, J.L. Peng, R.L. Greene, Nature 347 (1990) 369.
- [5] Z. Schlesinger, R.T. Collins, J.A. Caleise, D.G. Hinks, A.W. Mitchell, Y. Zheng, B. Dabrowski, N.E. Bickers, D.J. Scalapino, Phys. Rev. B 40 (1989) 6862.
- [6] H. Sato, T. Ido, S. Uchida, S. Tajima, M. Yoshida, K. Tanabe, K. Tatsuhara, N. Miura, Phys. Rev. B 48 (1993) 6617
- [7] F. Sharifi, A. Pargellis, R.C. Dynes, B. Miller, E.S. Hellman, J. Rosamilia, E.J. Hartford Jr., Phys. Rev. B 44 (1991) 12521.
- [8] R.J. Cava, B. Batlogg, J.J. Krajewski, R. Farrow, L.W. Rupp Jr., A.E. White, K. Short, W.F. Peck Jr., T. Kometani, Nature 332 (1988) 814.
- [9] W. Schmidbauer, A.J.S. Chowdbury, F. Wondre, M. Wanklyn, J.W. Hodby, Physica C 235–240 (1994) 759.
- [10] A. Carrington, R.W. Giannetta, J.T. Kim, J. Giapintzakis, Phys. Rev. B 59 (1999) 14173.
- [11] R. Prozorov, R.W. Giannetta, A. Carrington, F.M. Araujo-Moreira, Phys. Rev. B 62 (2000) 115.
- [12] H.T. Kim, K.Y. Kang, B.J. Kim, Y.C. Kim, W. Schmidbauer, J.W. Hodby, Physica C 341–348 (2000) 729.
- [13] J.F. Annet, N. Goldenfeld, A.J. Leggett, J. Low Temp. Phys. 105 (1996) 473.
- [14] P.J. Hirschfeld, N. Goldenfeld, Phys. Rev. B 48 (1993) 4219.